

# MoM scattering analysis of dihedral corner reflector: TALGAT verification

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**Abstract**—This paper presents the results of evaluating the effectiveness of using TALGAT system that based on the Method of Moments with pulse basis function for modelling dihedral corner reflector using the wire grid approach. TALGAT backscattering and radar cross section results are compared with those obtained experimentally and numerically using other methods published in other papers. Comparisons showed high consistency, confirming the effectiveness of using TALGAT for solving such problems, as well as for generating novel sparse wire grid scatterers in the near future.

**Keywords**—scatterer, TALGAT, method of moments, wire grid, dihedral corner reflector.

## I. INTRODUCTION

Researchers have long been interested in studying electromagnetic scattering from various structures. The development of advanced technology paved the way to the analysis of complex scattering structures. These structures are used in many fields today, including civilian [1] and military [2] applications. One of such structures that have received a significant attention are the corner reflectors [3, 4]. They have many advantages, including: high backscattering cross section (BSCS) over a wide angular range; operation without electrical power; mechanical simplicity in design; and ability to operate in harsh environments [5]. They are used in synthetic aperture radars [6], for ground target camouflage [7], and for target calibration in radar technology [8]. They also have many different shapes depending on the number of their scattering surfaces: dihedral, trihedral and octahedral etc. [9]. In this paper the dihedral corner reflector (DCR) that consists of two flat plates is considered.

In recent years, several studies focused on the analysis of perfectly conducting DCR [10, 11]. To analyze this structure, physical (PO) and geometrical (GO) optics individually or combined (PO\_GO), and physical theory of diffraction (PTD) are mainly used. The effectiveness of using these methods was demonstrated by good agreement with the experimental results [12]. The special feature of these methods when analyzing corner reflectors is the necessity to take into account the number of incident wave reflections from the reflector surfaces. This number depends on the angle between the two DCR surfaces. When the DCR angle is greater than  $90^\circ$  only singly and doubly reflected rays contribute to the scattered field. When this angle is between  $60^\circ$  and  $90^\circ$  only singly, doubly, and triply reflected rays contribute to the scattered field. While for angles less than

$60^\circ$ , it is necessary to consider rays undergoing more than three reflections [13]. This increases the complexity of analyzing DCR structures.

The method of moments (MoM) is widely used in modeling various electromagnetic scattering structures such as arbitrarily shaped ones [14], wire grid (WG) [15], etc. The core process of MoM is based on converting the electric field integral equations that describe the structure under study, its boundary conditions and the excitation of it into a system of linear equations. Then through the inversion of the impedance matrix solve this system using a computer [16]. This results in the current distribution over the surface of the considered structure, which makes it possible to determine the radiated or scattered field from it. This method allows one to analyze DCR without taking into account the number of wave reflections. Due to its nature, MoM considers the mutual influence between the structure elements in its impedance matrix, while the incident wave is considered as the excitation source of these elements. MoM is a method with simple algorithm and requires less resources than other methods. It can be used with different basis functions, one of which is the pulse basis functions, that have proved their effectiveness in analyzing complex WG scattering structures [17].

TALGAT is one of the currently known MoM-based systems that can be used to solve electromagnetic problems [18]. Recently, a MoM-based scattering module with pulse basis functions has been integrated into it to analyze various WG structures. The results obtained by this module were verified for simple wire structures. However, the effectiveness of using this module needs to be evaluated by verifying its results for different structures. Moreover, it is planned to be used for generating new sparse WG scatterers by employing the recently developed optimal current grid approximation approach [19]. This increases the relevance of evaluating its performance in analyzing various WG scatterers, of which DCRs are considered to be important ones. Therefore, the aim of this paper is to verify the results of modeling DCR scatterers as WGs obtained using TALGAT system. These results are compared with those obtained numerically using PTD [20], PO\_GO [21], PO [22], MoM with piecewise-sinusoidal (PWS) basis functions [23] and experimentally [20, 22, 23].

This paper is organized as follows: Section II provides the parameters of the incident plane waves, the considered DCR scatterers, and their equivalent WGs. Section III presents the BSCS and the radar cross section (RCS) reduction results of the considered scatterers using TALGAT and compares them with those published elsewhere to

evaluate their accuracy. Section IV summarizes the study conclusions and identifies directions for future research.

## II. STRUCTURES UNDER STUDY

To verify the results of MoM-based TALGAT with pulse basis functions obtained for DCR scatterers, structures of different sizes are considered. The isometric view of a solid perfectly conducting DCR structure is shown in Fig. 1a, while its equivalent WG structure is illustrated in Fig. 1b. DCR structure is formed by two rectangular plates (A and B) whose intersection coincides with the Oz axis. Both plates have a length of  $h$ , widths of  $w_1$  and  $w_2$  respectively. They are also symmetric with respect to the xOz plane and the angle between them is  $2\gamma$ . When modeling the considered DCR structures by WG, their plates are divided into cells of equal size and edge length of  $\Delta$  (Fig. 1b) (each WG cell edge is considered to be as one wire represented by one segment with a length equal to  $\Delta$ ). The value of  $\Delta$  is determined as:  $\lambda/6 \geq \Delta \geq \lambda/20$  [1], while the wire radius  $a$  is obtained by:  $a = \Delta/2\pi$ . To excite the scatterers, incident plane waves are used with linear vertical polarization ( $\theta$ -polarization) and directions determined by  $\varphi^{\text{inc}}$  and  $\theta^{\text{inc}}$ . The scattered waves have directions defined by  $\varphi^{\text{s}}$  and  $\theta^{\text{s}}$ , and are taken as  $\varphi^{\text{s}} = \varphi^{\text{inc}}$  and  $\theta^{\text{s}} = \theta^{\text{inc}}$  when calculating BSCS. In this paper, we considered 6 DCR structures differing in size and angle. Their parameters and the directions of the used incident plane wave to excite them are listed in Table 1. This table also reports the methods used to analyze these structures according to the papers that considered them.

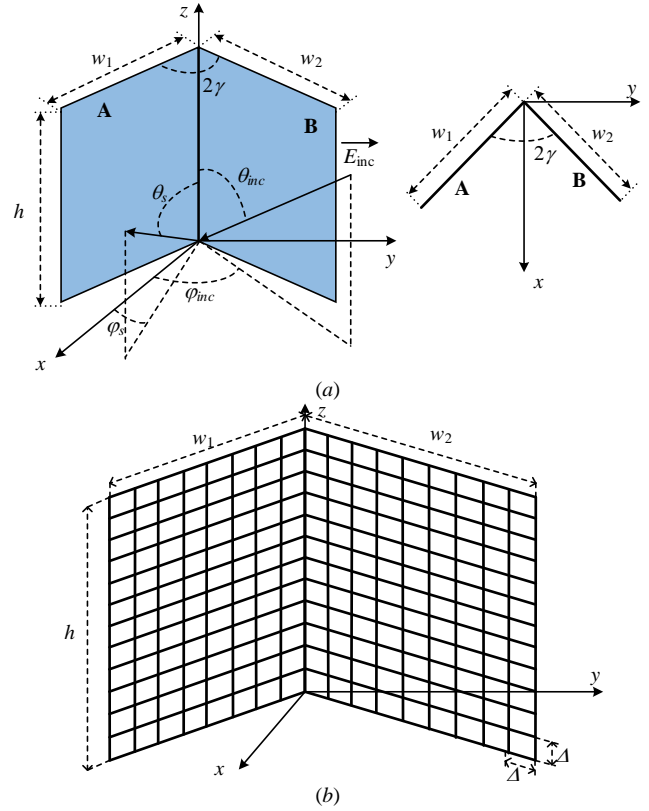


Fig. 1. DCR isometric view (a) and its equivalent WG (b).

TABLE I. ANALYSIS PARAMETERS OF THE CONSIDERED STRUCTURES

Structure	$w_1$ , m	$w_2$ , m	$h$ , m	$a$ , m	Cell numbers	$\varphi^{\text{inc}}, ^\circ$	$\theta^{\text{inc}}, ^\circ$	$f$ , GHz	$2\gamma, ^\circ$	Analysis
$S_1$ [20]	0.18	0.18	0.18	0.0009	33×33×33	0–360	90	9.4	90	PTD / measurement
									98	
									77	
$S_2$ [21]	7.16	4.77	4.77	0.026	43×29×29	0–360	90	0.3	90	PO+GO
$S_3$ [22]	2.5	2.5	5	0.019	20×20×40	0–360	90	0.3	85	PO / measurement
							85		95	
$S_4$ [22]	5	5	10	0.027	30×30×60	0–360	90	0.3	90	PO
									60	
$S_5$ [22]	6	6	10	0.024	40×40×66	0–360	90	0.3	97	PO / measured
									97	
$S_6$ [23]	0.5	0.5	1	0.008	10×10×20	0–360	90	0.3	130	MoM-PWS / measurement
						0	0–360			

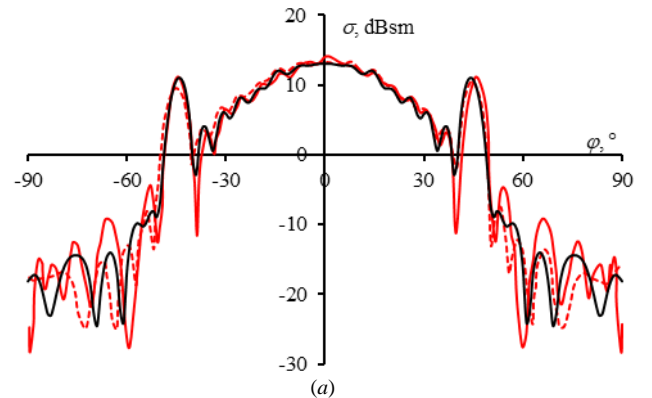
## III. RESULTS

First, BSCS for  $S_1$  with different  $2\gamma$  angles are considered. Their values obtained in TALGAT are compared with those obtained experimentally and numerically using PTD in [20] (Fig. 2). It is demonstrated that TALGAT results are in an acceptable agreement with PTD results in the main lobe, but the difference between them increases in the side lobe, especially when  $2\gamma=98$  and  $77^\circ$ . These deviations are summarized in Table 2. It can be seen from Fig. 2a that the BSCS for DCR scatterer with  $2\gamma=90^\circ$  reach their maximum values at  $\varphi^{\text{inc}}=0^\circ$ . They slightly change over the azimuth range of  $(-15^\circ; 15^\circ)$ , but rapidly decrease beyond it. This is due to the fact that in this range most of the scattered energy returns in the opposite direction of the incident wave one.

TABLE II. BSCS DEVIATIONS FROM [20] FOR  $S_1$

$2\gamma, ^\circ$	Deviations in $\varphi^{\text{s}}=0^\circ, \theta^{\text{s}}=90^\circ$ , dB		Maximum deviations, dB	
	PTD	measured	PTD	measured
90	1.3	0	10	9
98	0.2	4.8	8	6
77	0.8	1.1	8	3.8

In addition, BSCS values increase suddenly in the ranges of  $(-50^\circ; -40^\circ)$  and  $(40^\circ; 50^\circ)$ . This can be explained by the fact that the plane wave had a direction almost perpendicular to one of the DCR plates. On the other hand, when  $2\gamma=98^\circ$  or  $77^\circ$ , the scattered energy returns in a direction that differs from the incident wave one in the main lobe region (Fig. 2b and 2c).



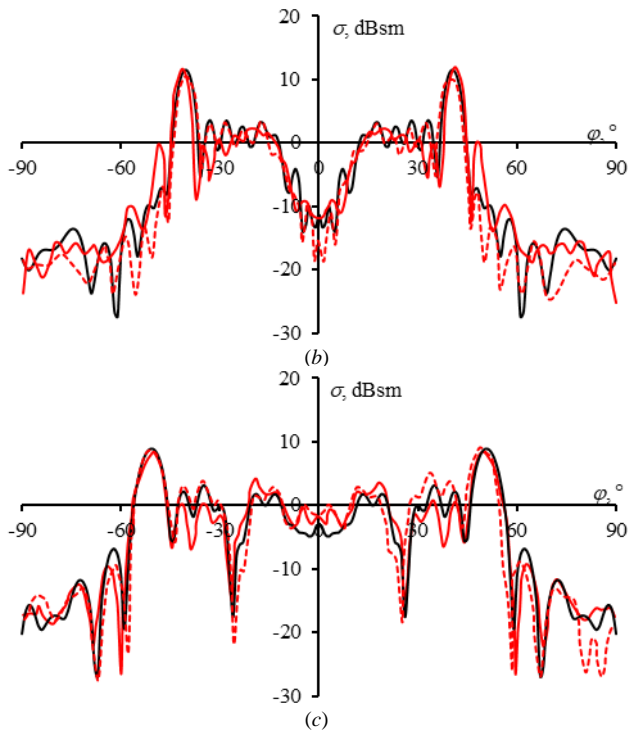


Fig. 2. BSCS for  $S_1$  when  $2\gamma=90^\circ$  (a),  $98^\circ$  (b),  $77^\circ$  (c) obtained using TALGAT (—), PTD (—), and experimentally (---) [20].

Next, the BSCS for  $S_2$  obtained using TALGAT were compared with those calculated using PO\_GO in [21] (Fig. 3). It can be seen that TALGAT results are also in good agreement with those obtained PO\_GO (their highest values differ maximumly by about 2 dB). Fig. 3 also shows that the scattered field maximums shifted by about  $10^\circ$  toward the side of the DCR plate with bigger width. The side lobe in the range of  $\varphi_{inc}=(40; 50)$  has BSCS maximum value (27.5 dB) greater than that in  $(-50; -40)$  (24.5 dB). This can be explained by the fact that plate A has a larger size than plate B, thus the scattered field from it will have a larger magnitude.

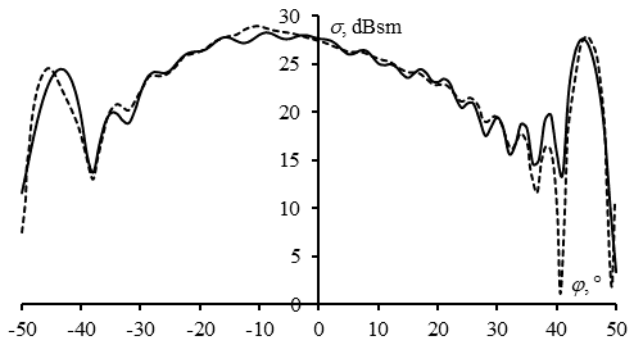


Fig. 3. BSCS for  $S_2$  obtained using TALGAT (—) and PO\_GO (---) [21].

After that, the RCS reductions for  $S_3$  when  $2\gamma=85^\circ$  obtained using TALGAT were compared with those calculated using PO and measured in [22] (Fig. 4). These reductions are determined by the difference between the RCS values for DCRs with an angle other than  $90^\circ$  and their values when it is equal to it [22]. It can be seen that TALGAT results are also in good agreement with those obtained experimentally (maximum deviation is 0.8 dB). This difference is more noticeable compared to PO (maximum deviation is 1.8 dB). Nevertheless, TALGAT results agree with measurements better than PO ones.

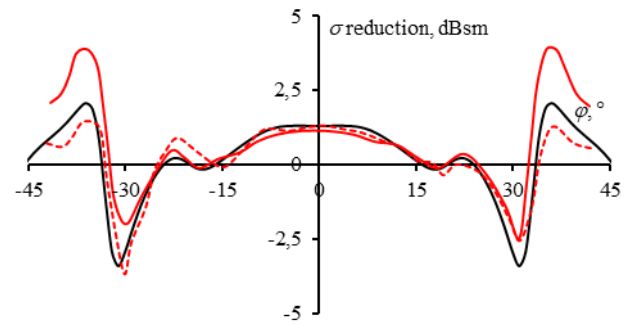


Fig. 4. RCS reductions for  $S_3$  when  $2\gamma=85^\circ$  obtained using TALGAT (—), PO (—), and experimentally (---) [22].

Next, the BSCS for  $S_3$  when  $2\gamma=95^\circ$  at  $\theta^{inc}=90^\circ$  and  $\theta^{inc}=85^\circ$  obtained using TALGAT were compared with those calculated using PO and measured in [22] (Fig. 5). It can be seen that TALGAT results are also in good agreement with those obtained numerically using PO (maximum deviation is less than 1.9 dB) and experimentally (maximum deviation is less than 2 dB). It can be noticed that the scattering field at  $\theta^{inc}=90^\circ$  is larger than at  $\theta^{inc}=85^\circ$ . This can be explained by the fact that at  $\theta^{inc}=85^\circ$  the maximum scattering energy will be directed to  $\theta=95^\circ$ , not in the opposite direction of the excitation wave ( $\theta=85^\circ$ ).

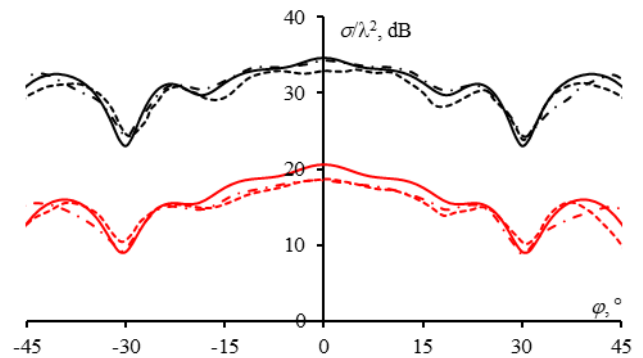


Fig. 5. BSCS for  $S_3$  when  $2\gamma=95^\circ$  at  $\theta_{inc}=90^\circ$  (black) and  $\theta_{inc}=85^\circ$  (red) obtained using TALGAT (—), PO (---) and experimentally (---) [22].

Then, the calculated BSCS for  $S_4$  when  $2\gamma=90^\circ$  and  $60^\circ$  using TALGAT are compared with those obtained using PO in [22] (Fig. 6). When  $2\gamma=90^\circ$ , it can be seen that TALGAT results are also in good agreement with those obtained PO (with deviation about 1 dB in  $\varphi^{inc}=0^\circ$ , and maximum deviation about 3 dB). While the difference between the compared results is bigger when  $2\gamma=60^\circ$  (with deviation about 1.3 dB in  $\varphi^{inc}=0^\circ$ , and maximum deviation about 13 dB).

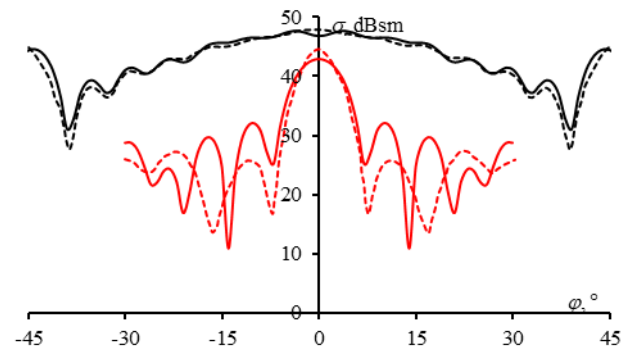


Fig. 6. BSCS for  $S_4$  when  $2\gamma=90^\circ$  (black lines) and  $2\gamma=60^\circ$  (red lines) obtained using TALGAT (—) and PO (---) [22].

Moreover, the calculated BSCS for  $S_5$  using TALGAT are compared with those obtained using PO and experimentally from [22] (Fig. 7). It can be seen that TALGAT results agree well with the measured results (with deviation about 1.8 dB in  $\varphi=0^\circ$ , and maximum deviation about 8.6 dB). However, a larger deviation is observed by comparing them to those obtained numerically using PO (with deviation about 8.4 dB in  $\varphi=0^\circ$ , and maximum deviation about 16 dB).

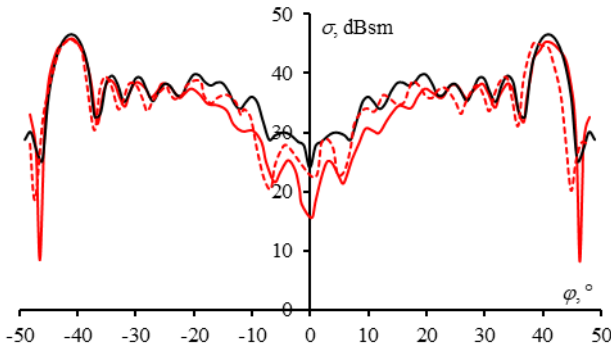


Fig. 7. BSCS for  $S_5$  obtained using TALGAT (—), PO (—) and experimentally (---) [22].

Finally, the calculated BSCS results for  $S_6$  in the  $\theta=90^\circ$  and  $\varphi=0^\circ$  planes using TALGAT are compared with those obtained experimentally and numerically using MoM with PWS in [23] (Fig. 8). It can be seen that TALGAT results are also in good agreement with those obtained experimentally (the maximum deviation is less than 1.5 dB in the  $\theta=90^\circ$  plane and less than 0.25dB – in the  $\varphi=0^\circ$  plane). Fig. 8 shows also that TALGAT results are closer to the measured ones than those of MoM with PWS. The maximum deviations calculated when comparing TALGAT and MoM with PWS results are about 9 dB in the  $\theta=90^\circ$  plane and 3.5 dB – in the  $\varphi=0^\circ$  plane.

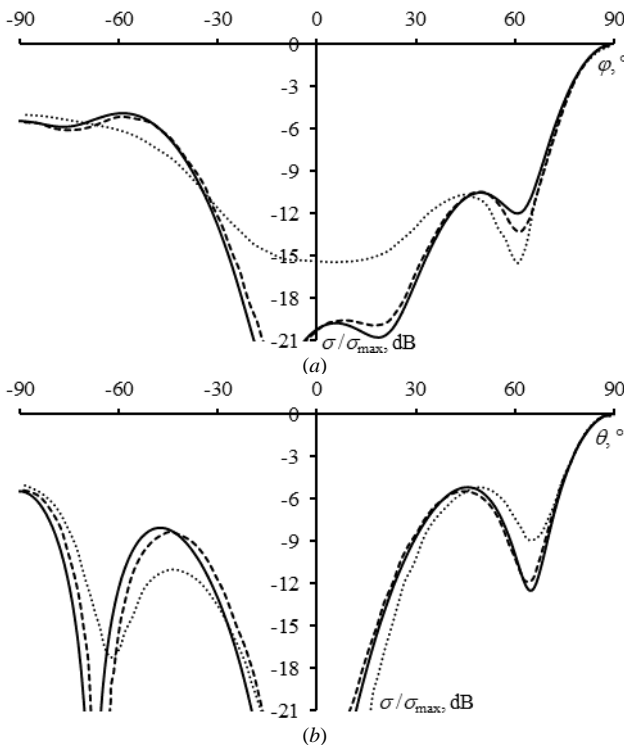


Fig. 8. BSCS for  $S_6$  in the  $\theta=90^\circ$  (a) and  $\varphi=0^\circ$  (b) planes obtained using TALGAT (—), MoM with PWS (····) and experimentally (---) [23].

In general, the results of modeling DCR scatterers using MoM-based WG with pulse basis functions in TALGAT system are quite similar to those obtained numerically using other methods and experimentally. Based on that, it is evident to conclude that TALGAT can be efficiently used to analyze complex scattering structures.

#### IV. CONCLUSION

The BSCS and RCS reduction results for DCR scatterers obtained in TALGAT are verified by detailed comparisons with those obtained numerically using PTD, PO, PO\_GO and experimentally in other works. These comparisons yielded a high degree of consistency, proving the accuracy and effectiveness of MoM-based TALGAT with pulse basis functions in solving such problems. This study paves the way to explore the potential of TALGAT in analyzing a wider range of complex WG scatterers. Notably, the successful verification on DCR scatterers forms the foundation for future investigations on novel sparse WG scatterers.

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